

Regional Cooperation for
Limited Area Modeling in Central Europe



ACC  RD

A Consortium for CONvection-scale modelling
Research and Development

Dynamics in LACE - towards hectometric scales

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Czech
Hydrometeorological
Institute



HungaroMet



ARSO METEO
Slovenia

- ❑ **Dynamical core in ACCORD**
- ❑ **SI time scheme**
- ❑ **Orographic terms in linear model (based on work of Jozef Vivoda and Fabrice Voitus)**
 - ❑ **Some equations**
 - ❑ **Idealised tests**
 - ❑ **Real simulations @200m**
- ❑ **Vertical velocity definition in nonlinear model (based on work of Fabrice Voitus)**
 - ❑ **Real simulations @200m**

Basic equations

- hydrostatic primitive equation system (HPE) or Euler equations (EE); recently implemented quasi elastic equation system (QE)
- prognostic variables $\vec{v}, T, q_s = \ln(\pi_s)$, in EE with $w, \hat{q} = \ln(\frac{p}{\pi})$

Discretization

- spectral transform method for horizontal direction
- hybrid vertical coordinate η based on hydrostatic pressure $\pi(\eta) = A(\eta) + B(\eta)\pi_s$;
 $A(top) = B(top) = 0, A(bottom) = 0, B(bottom) = 1$
- finite differences or finite elements for vertical direction discretization
- semi-implicit or iterative centred implicit scheme for time discretization
- semi-Lagrangian advection

System evolution

$$\frac{dX}{dt} = \mathcal{M}X$$

Using linear model \mathcal{L}^* we get

$$\frac{dX}{dt} = \mathcal{L}^* \overline{[X]}^t + (\mathcal{M} - \mathcal{L}^*)X$$

and discretize in time to obtain

Linearization

$$X = \mathbf{X}^* + \mathbf{X}', \quad \mathcal{M} \rightarrow \mathcal{L}^*$$

Semi-implicit scheme

$$\frac{\mathbf{X}^+ - \mathbf{X}^0}{\Delta t} = \mathcal{L}^* \left(\frac{\mathbf{X}^+ + \mathbf{X}^0}{2} \right) + (\mathcal{M} - \mathcal{L}^*)\mathbf{X}^{+\frac{1}{2}}$$

or

Iterative centered implicit scheme

$$\frac{\mathbf{X}^{+(n)} - \mathbf{X}^0}{\Delta t} = \frac{\mathcal{L}^* \mathbf{X}^{+(n)} + \mathcal{L}^* \mathbf{X}^0}{2} + \frac{(\mathcal{M} - \mathcal{L}^*)\mathbf{X}^{+(n-1)} + (\mathcal{M} - \mathcal{L}^*)\mathbf{X}^0}{2}$$

We know that both can be second order accurate in time when some care is taken (averaging along semi-Lagrangian trajectory).

Temperature

$$\frac{dT}{dt} = \frac{\kappa T}{\kappa - 1} (D + d)$$

Horizontal momentum

$$\frac{d\vec{v}}{dt} = -RT \frac{\nabla \pi}{\pi} - \nabla \phi - RT \nabla \hat{q} - \frac{1}{m} \frac{\partial(p - \pi)}{\partial \eta} \nabla \phi$$

Vertical momentum

$$\frac{dw}{dt} = \frac{g}{m} \frac{\partial(p - \pi)}{\partial \eta}$$

Pressure departure

$$\frac{d\hat{q}}{dt} = \frac{1}{\kappa - 1} (D + d) - \frac{1}{\pi} \frac{d\pi}{dt}$$

Surface pressure

$$\frac{dq_s}{dt} = -\frac{1}{\pi_s} \int_0^1 \nabla \cdot (m\vec{v}) d\eta$$

Diagnostic relations

$$\frac{d\pi}{dt} = \vec{v} \cdot \nabla \pi - \int_0^\eta \nabla \cdot (m\vec{v}) d\eta'$$

$$\phi = \phi_s - \int_\eta^1 \frac{mRT}{p} d\eta'$$

$$d = -\frac{p}{mRT} \left(\nabla \phi \frac{\partial \vec{v}}{\partial \eta} - g \frac{\partial w}{\partial \eta} \right)$$

Definitions

$$D = \nabla \cdot \vec{v}$$

$$\kappa = \frac{c_p}{R}$$

$$m = \frac{\partial \pi}{\partial \eta}$$

Temperature

$$\frac{dT}{dt} = \frac{\kappa T}{\kappa - 1} (D + d)$$

Horizontal momentum

$$\frac{d\vec{v}}{dt} = -RT \frac{\nabla \pi}{\pi} - \nabla \phi - RT \nabla \hat{q} - \frac{1}{m} \frac{\partial(p - \pi)}{\partial \eta} \nabla \phi$$

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Definitions

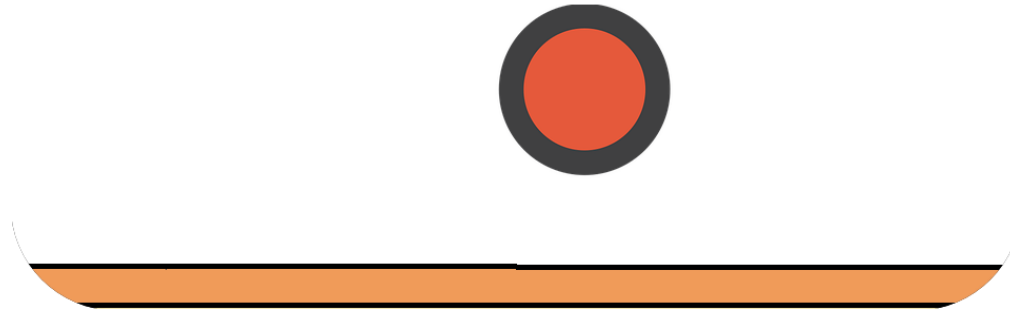
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$$\kappa = \frac{c_p}{R}$$

$$m = \frac{\partial \pi}{\partial \eta}$$

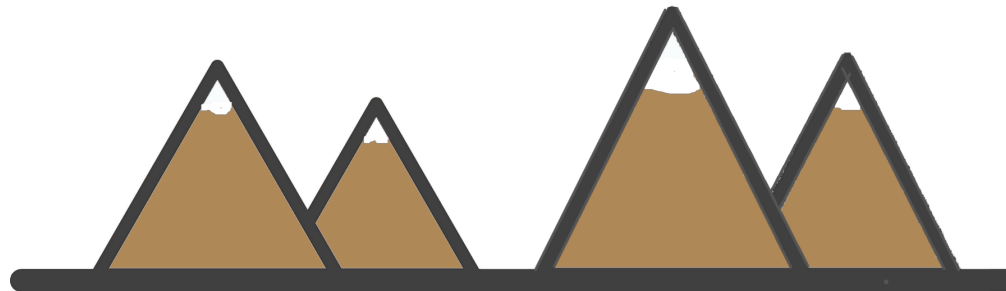
Current state:

- stationary
- resting
- hydrostatically balanced (π_s^*)
- dry
- isothermal (T^*)
- with constant orography ($\nabla\phi^* = 0$)



New:

- stationary
- resting
- hydrostatically balanced (π_s^*)
- dry
- isothermal (T^*)
- with constant orographic slope (in absolute value, $|\nabla\phi^*| \neq 0$)



Temperature

$$\frac{\partial T}{\partial t} = \frac{\kappa T^*}{\kappa - 1} (D + d)$$

Horizontal momentum

$$\frac{\partial \vec{v}}{\partial t} = -RT^* \frac{\nabla \pi}{\pi^*} - \nabla \phi - RT^* \nabla \hat{q} - \frac{1}{m^*} \frac{\partial \pi^* \hat{q}}{\partial \eta} \nabla \phi^*$$

Vertical momentum

$$\frac{\partial w}{\partial t} = \frac{g}{m^*} \frac{\partial \pi^* \hat{q}}{\partial \eta}$$

Pressure departure

$$\frac{\partial \hat{q}}{\partial t} = \frac{1}{\kappa - 1} (D + d) + \frac{1}{\pi^*} \int_0^\eta m^* D d\eta'$$

Surface pressure

$$\frac{\partial q_s}{\partial t} = -\frac{1}{\pi_s^*} \int_0^1 m^* D d\eta$$

Diagnostic relations

$$\begin{aligned} \nabla \phi &= \nabla \phi_s - \int_\eta^1 \nabla \left(\frac{mRT}{p} \right) d\eta' \\ \nabla \phi^* &= g\Lambda^* S^*(\eta) \\ d &= -\frac{p}{mRT} \left(\nabla \phi \frac{\partial \vec{v}}{\partial \eta} - g \frac{\partial w}{\partial \eta} \right) \end{aligned}$$

Definitions

$$\begin{aligned} \Lambda^* &= \frac{1}{g} \|\nabla \phi_s\|^* \\ S^*(\eta) &= \frac{B(\eta) \pi_s^*}{\pi^*(\eta)} \\ m^* &= \frac{\partial \pi^*}{\partial \eta} \end{aligned}$$

Modified vertical divergence

$$d = \frac{p}{mRT} \left(\nabla \phi \frac{\partial \vec{v}}{\partial \eta} - g \frac{\partial w}{\partial \eta} \right)$$

Time evolution in linear model

$$\frac{\partial \vec{v}}{\partial t} = \mathbb{A} - \Lambda^* S^* (\eta) \mathbb{B}$$

$$\frac{\partial w}{\partial t} = \mathbb{B}$$

$$\begin{aligned} \frac{\partial d}{\partial t} &= \frac{1}{RT^*} \left[\nabla \phi^* \partial^* \left(\frac{\partial \vec{v}}{\partial t} \right) - g \partial^* \left(\frac{\partial w}{\partial t} \right) \right] \\ &= \frac{1}{RT^*} \left[g \Lambda^* S^* (\eta) \partial^* \mathbb{A} - g \Lambda^{*2} S^* (\eta) (S^* \partial^* \mathbb{B} + \mathbb{B} \partial^* S^*) - g \partial^* \mathbb{B} \right] \end{aligned}$$

where

$$\partial^* X = \frac{\pi^*}{m^*} \frac{\partial X}{\partial \eta}$$

We omit the first order terms in Λ^* and then $\frac{\partial \vec{v}}{\partial t}$ is unchanged and all operators of the RHS of $\frac{\partial d}{\partial t}$ apply on \hat{q} .

Time evolution in linear model

$$\begin{aligned} \frac{\partial \vec{v}}{\partial t} &= \mathbb{A} - \cancel{\Lambda^* S^* (\eta)} \mathbb{B} \\ \frac{\partial d}{\partial t} &= \frac{1}{RT^*} \left[\cancel{g \Lambda^* S^* (\eta)} \partial^* \mathbb{A} - g \Lambda^{*2} S^* (\eta) (S^* \partial^* \mathbb{B} + \mathbb{B} \partial^* S^*) - g \partial^* \mathbb{B} \right] \end{aligned}$$

Finally, since $\mathbb{B} = g(\partial^* + 1)\hat{q}$

Time evolution in linear model

$$\begin{aligned} \frac{\partial \vec{v}}{\partial t} = \mathbb{A} &\rightsquigarrow \frac{\partial D}{\partial t} = \frac{\partial(\nabla \cdot \vec{v})}{\partial t} = \nabla \cdot \mathbb{A} \\ \frac{\partial d}{\partial t} &= \mathcal{L}_{new}^* \hat{q} \end{aligned}$$

We can define

New vertical Laplacian operator

$$\mathcal{L}_{new}^* = \alpha \partial^* (\partial^* + 1) + \beta (\partial^* + 1)$$

$$\alpha = 1 + \Lambda^{*2} S^{*2}(\eta)$$

$$\beta = \Lambda^{*2} S^*(\eta) \partial^* S^*(\eta)$$

$$\Lambda^* = 0 : \mathcal{L}_{new}^* \rightarrow \mathcal{L}_v^* = \partial^* (\partial^* + 1)$$

$$S^*(\eta) = 0 : \mathcal{L}_{new}^* \rightarrow \mathcal{L}_v^*$$

How to discretize the proposed solution?

New discretized vertical Laplacian operator

$$[\partial^* (\partial^* + 1) X]_l = \dots$$

$$[(\partial^* + 1) X]_l = \dots$$

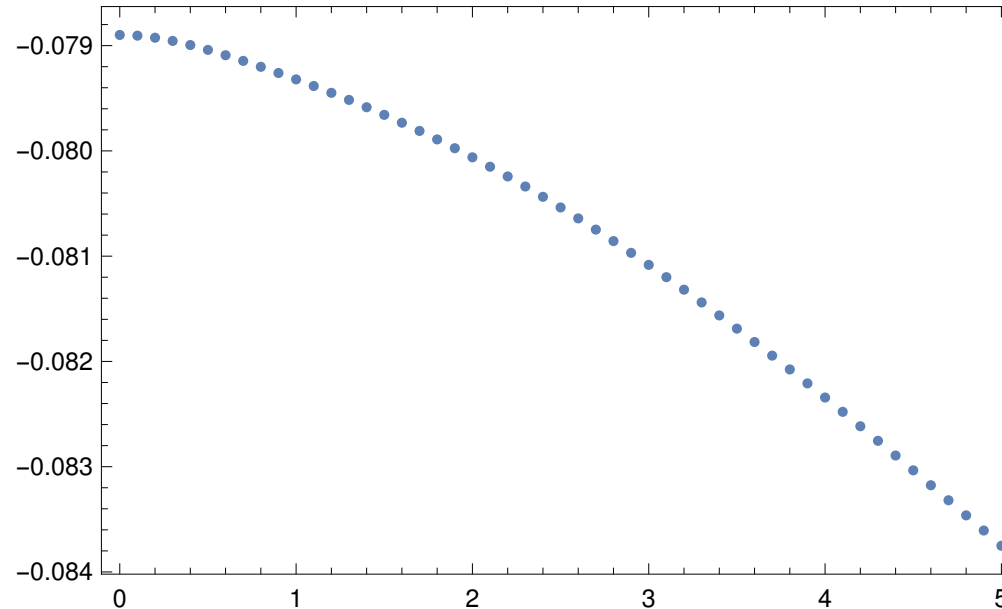
$$S^*(\eta) = \dots$$

$$\partial^* S^*(\eta) = \dots$$

How to set boundary conditions?

Does \mathcal{L}_{new}^* have only real and negative eigenvalues?

For an example of 87 vertical levels used in Czech operations we are safe.

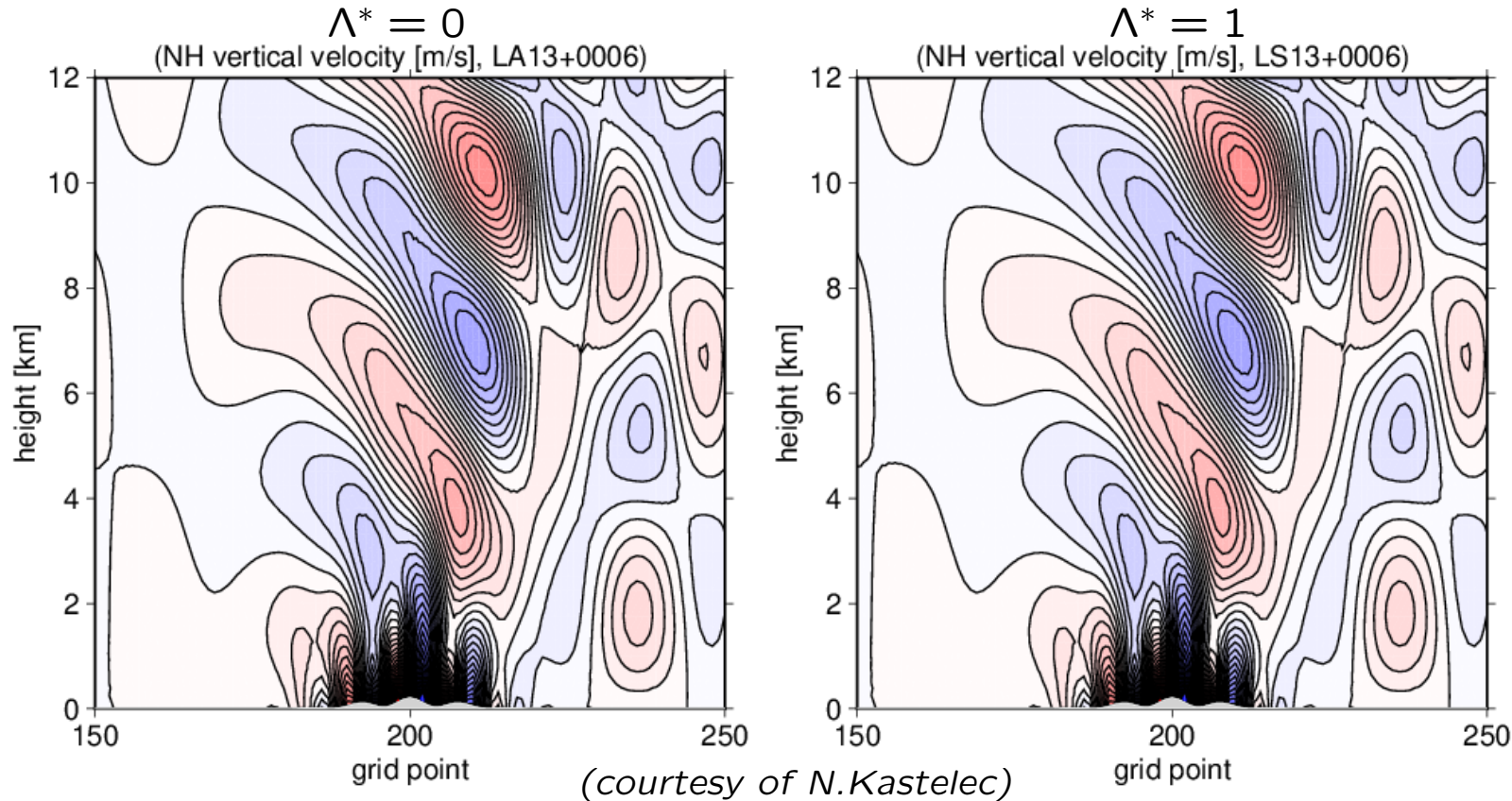


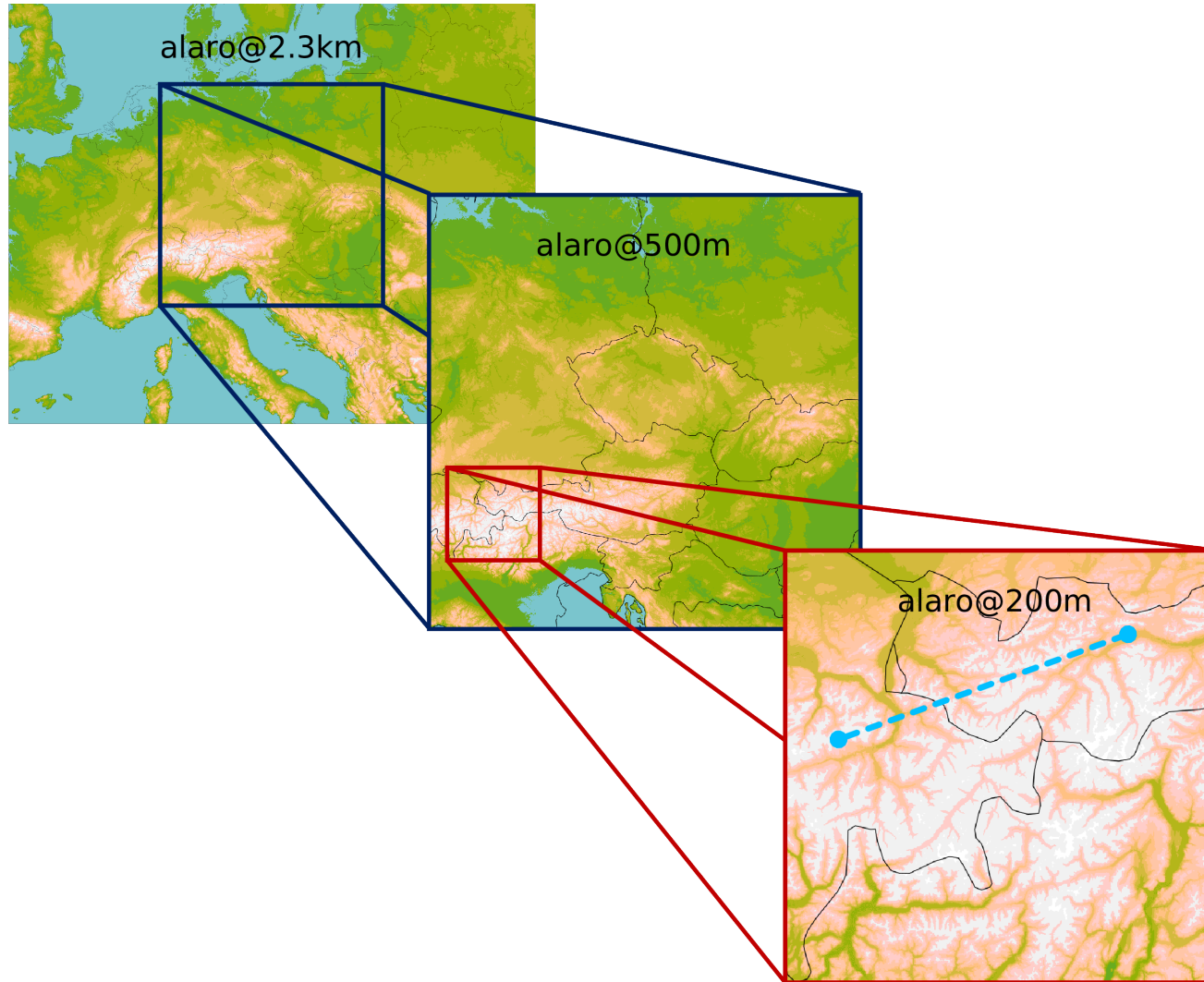
(courtesy of N.Kastelec)

Max value of eigen values depending on Λ

Then we can eliminate the discretized equations up to horizontal divergence D and solve the Helmholtz equation for D .

Vertical velocity for the Schär mountain case depending on Λ^* . ($\Delta x = 500\text{ m}$, $\Delta z = 250\text{ m}$, mountain height $h = 250\text{ m}$, $T_0 = 288\text{ K}$, $u_0 = 10\text{ m/s}$, $\Delta t = 32\text{ s}$)





← We show results of 200m simulation over Alps.

The basic algorithmic choices for ALARO configurations @200m are:

Dynamical core

- ❑ semi-Lagrangian advection scheme with 4 iterations for trajectory calculation
- ❑ PC time scheme with one iteration, cheap variant (SL trajectories are not recalculated in corrector)
- ❑ modified vertical divergence d4 for vertical motion, transformation to vertical velocity w in the non-linear model
- ❑ reference values of the linear model: SITR=300K, SITRA=100K, SIPR=900hPa
- ❑ no decentering
- ❑ semi-Lagrangian horizontal diffusion applied on all model variables + TKE, TTE, hydrometeors
- ❑ linear truncation for all spectral fields except orography; quadratic truncation of orography

ALARO physics

- ❑ radiation scheme ACRANE2
- ❑ turbulence and shallow convection scheme TOUCANS, model 2
- ❑ scale aware deep convection and microphysics scheme 3MT

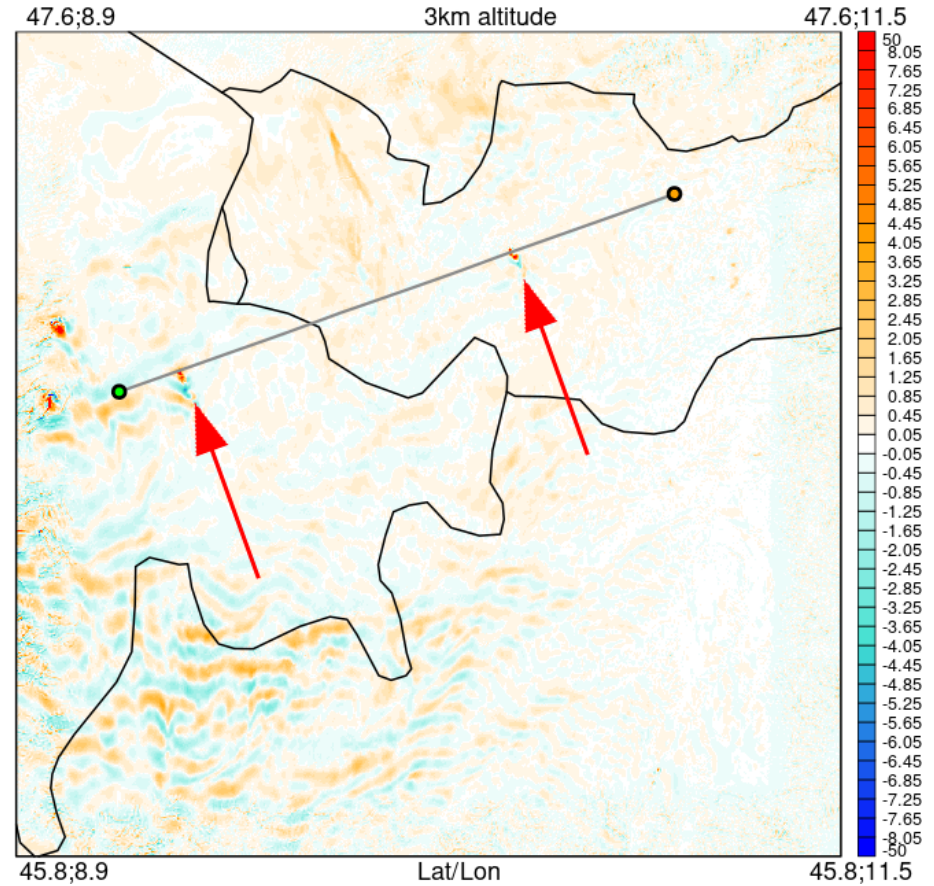
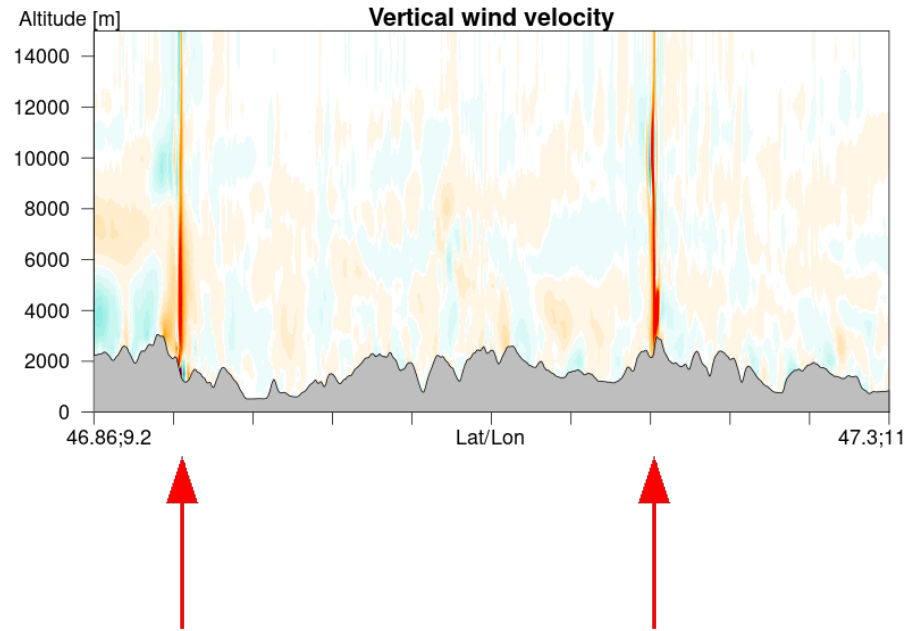
Initialization

- ❑ initialization with 3DVAR + surface DA (canari) for 2.325km run; dynamical adaptation + DFI for 500m and 200m runs

Particular choices for ALARO@200m:

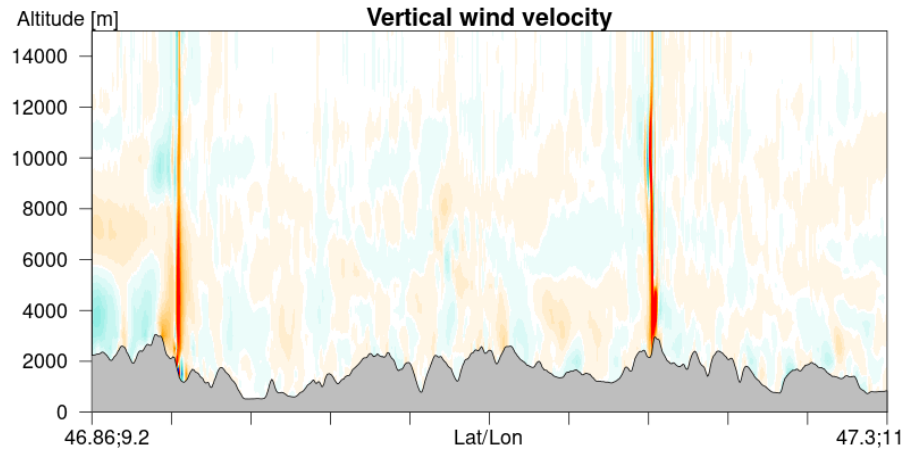
- ❑ cubic truncation of orography
- ❑ SITRA=50K
- ❑ no 3MT (deep convection), only STRAPRO (stratiform precipitation)

Vertical velocity for the alpine case 19 August 2022 OUTC + 24hours.



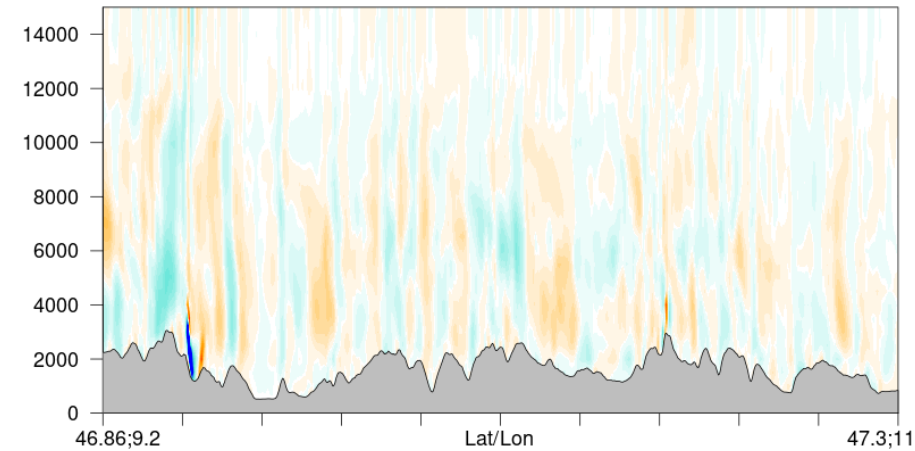
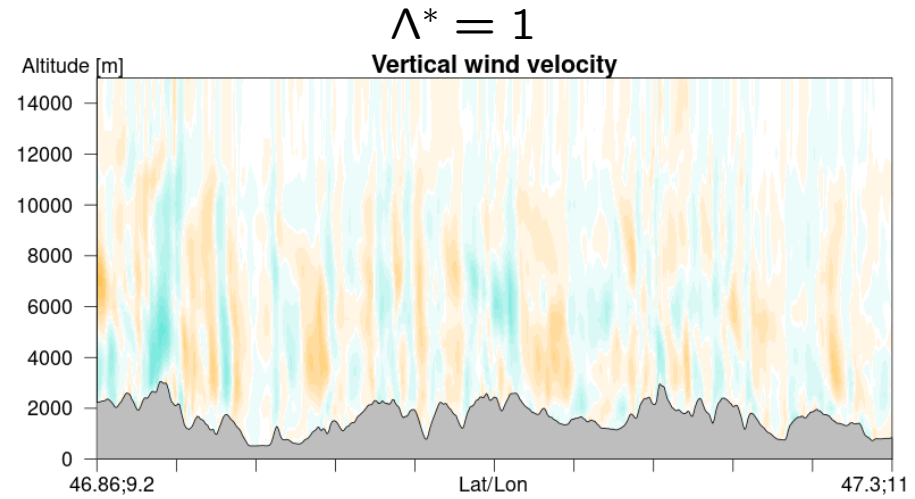
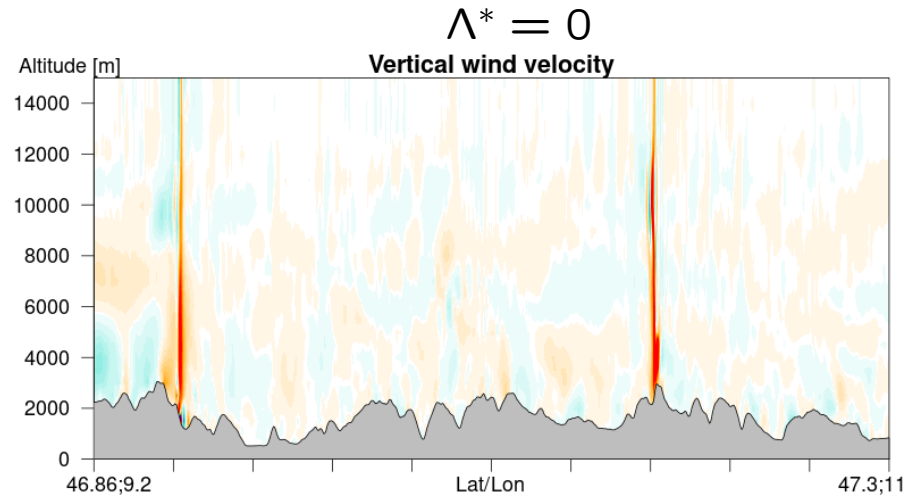
Vertical velocity depending on Λ^* for the alpine case 19 August 2022 OUTC + 24hours.

$$\Lambda^* = 0$$

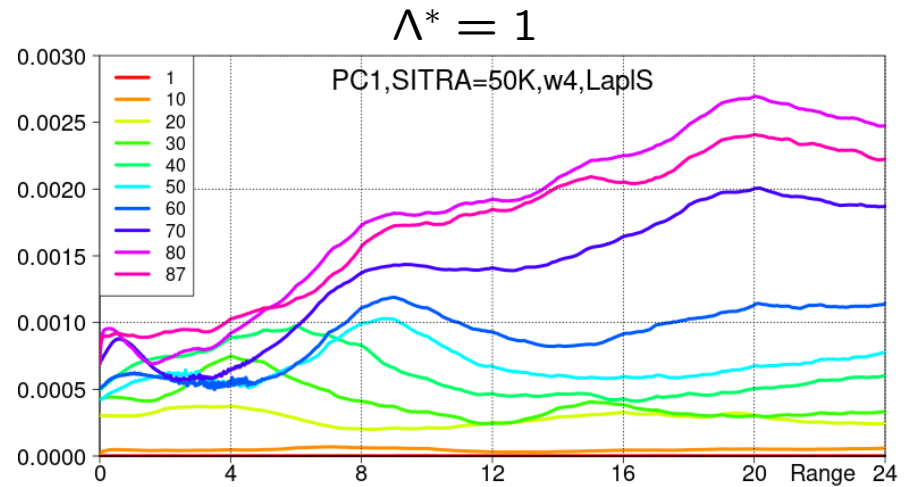
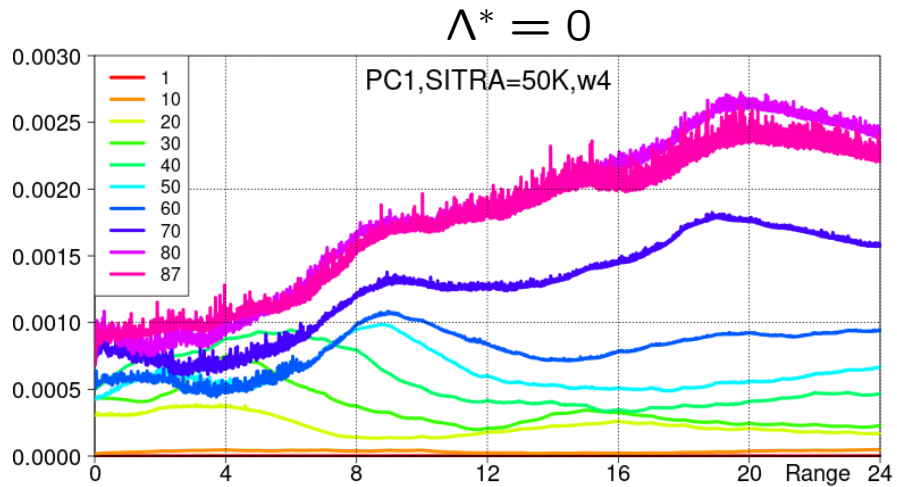


With additional iterations of the SI scheme, the integration crashes.

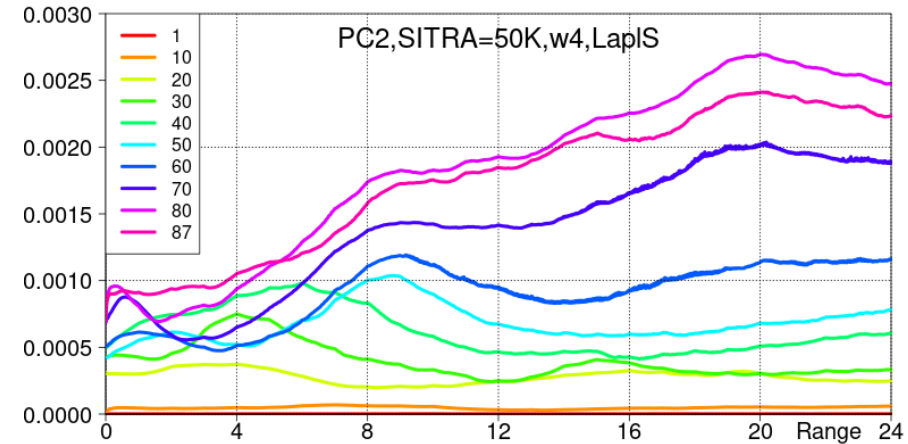
Vertical velocity depending on Λ^* for the alpine case 19 August 2022 OUTC + 24hours.



Averaged spectral norms of vertical divergence for the alpine case 19 August 2022 OUTC + 24hours.



PC+1



PC+2

Conclusions

- ❑ We must continue our efforts.
- ❑ We plan to test various possible definitions of vertical function $S^*(\eta)$.
- ❑ We plan to test various possible discretizations including boundary conditions.
- ❑ We plan to make further idealised tests and real simulations.
- ❑ If the time scheme allows a source of a noise, further time iterations may not help to stabilize the scheme. To the contrary, the scheme with further iterations may show even less stability.

Fabrice Voitus proposed a modification of the vertical velocity defined in the model to simplify the bottom boundary condition and allow more precise calculation.

<https://events.ecmwf.int/event/167/contributions/1379/attachments/794/1401/AS2020-Voitus.pdf>

Current state

$$w = \frac{dz}{dt}$$

$$d = -g \frac{p}{mRT} \frac{\partial w}{\partial \eta} + X$$

$$X = \frac{p}{mRT} \frac{\partial \mathbf{V}}{\partial \eta} \nabla \phi$$

$$\frac{\partial d}{\partial t} = -\frac{g^2}{RT^*} \mathcal{L}_v^* \hat{q}$$

$$\mathcal{L}_v^* = \partial^* (\partial^* + 1)$$

Fabrice's definition

$$W = w - \mathbf{V} \cdot \mathbf{S}(\eta) \frac{1}{g} \nabla \phi_s$$

$$d = -g \frac{p}{mRT} \frac{\partial W}{\partial \eta} + X$$

$$X = \frac{p}{mRT} \frac{\partial \mathbf{V}}{\partial \eta} \nabla (\phi - \mathbf{S}(\eta) \phi_s) - \frac{p}{mRT} \mathbf{V} \frac{\partial \nabla \phi}{\partial \eta}$$

$$\frac{\partial d}{\partial t} = -\frac{g^2}{RT^*} \mathcal{L}_{mod}^* \hat{q}$$

$$\mathcal{L}_{mod}^* = \partial^* (\partial^* + 1 + \Lambda^{*2} \mathbf{S}(\eta) \gamma^*)$$

Then discretization of these terms involves interpolations between half levels (where w is represented) and full levels (where \mathbf{V} is represented) and is cumbersome ...

Current state

in the non-linear model

$$w_s = \mathbf{V}_s \frac{1}{g} \nabla \phi_s$$
$$\frac{dw_s}{dt} = \frac{d\left(\mathbf{V}_s \frac{1}{g} \nabla \phi_s\right)}{dt}$$

while in the linear model explicit guess of d is calculated consistently and used in the implicit part

Complicated and not exact! It is a source of noise which may grow.

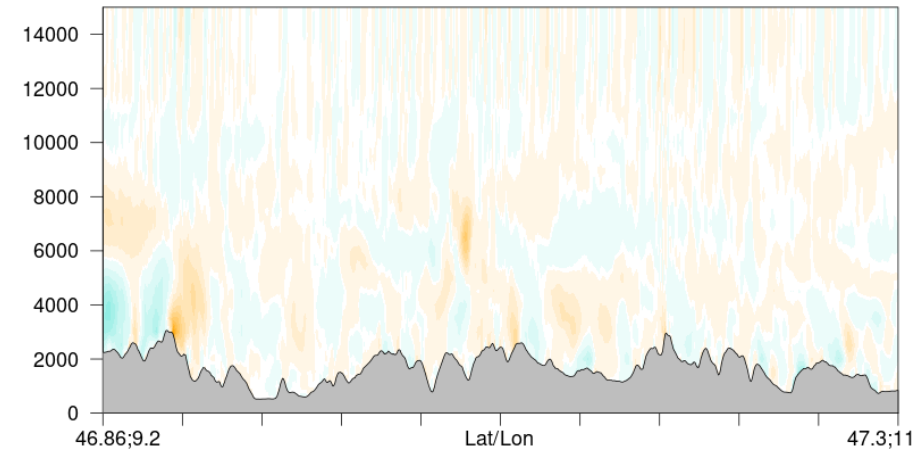
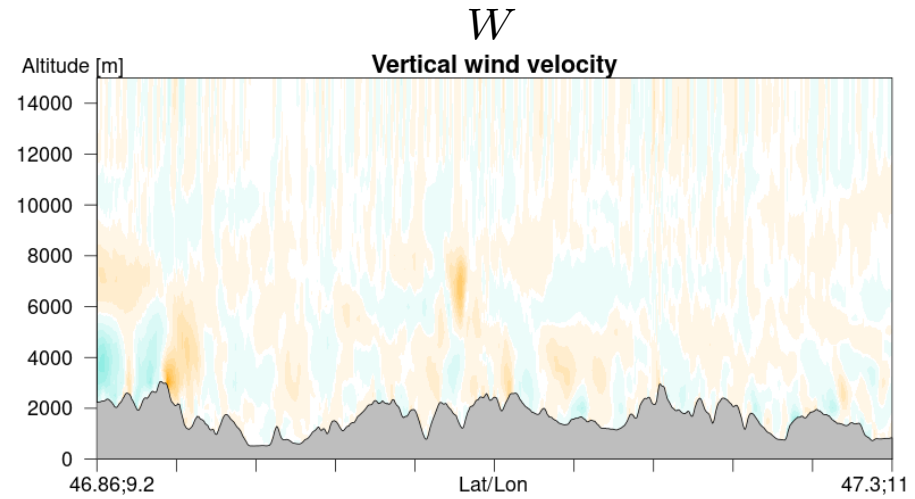
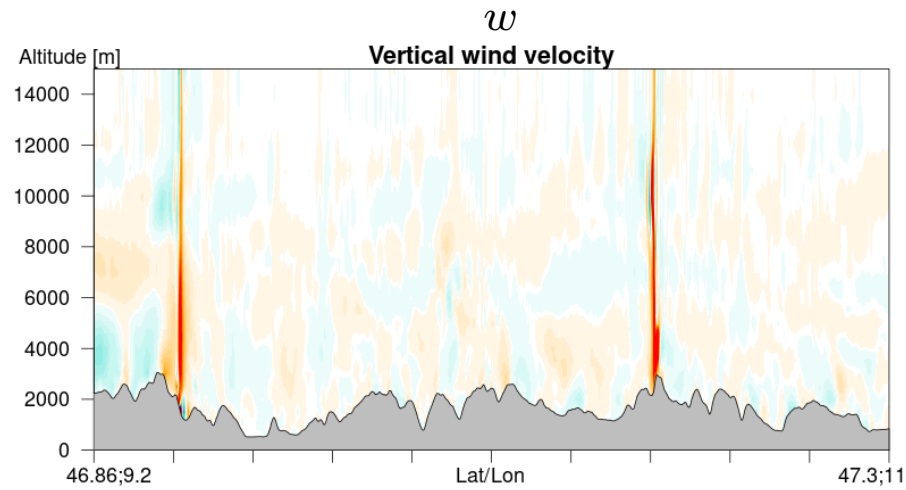
Fabrice's definition

in the linear and non-linear model

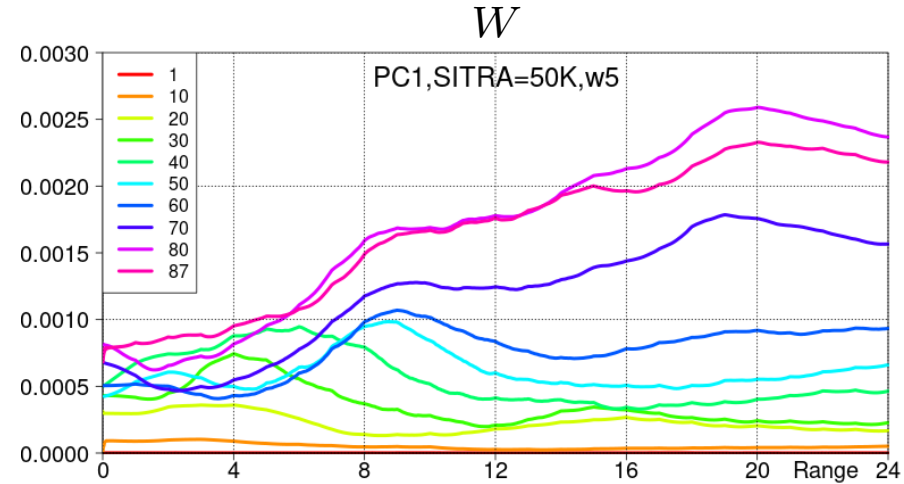
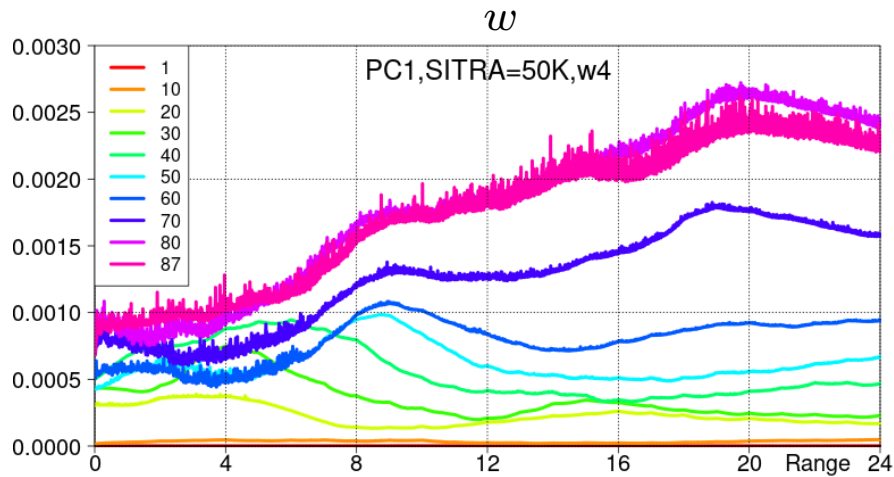
$$W_s = 0$$
$$\frac{dW_s}{dt} = 0$$

Easily applicable!

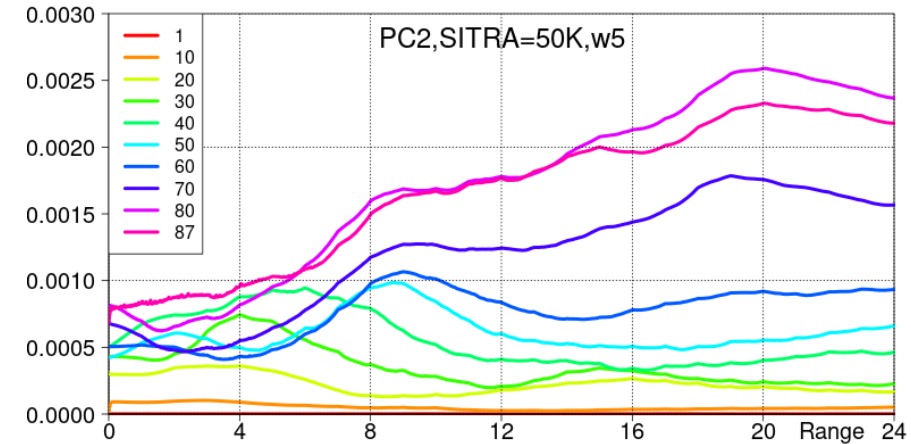
Vertical velocity for the alpine case 19 August 2022 OUTC + 24hours.



Averaged spectral norms of vertical divergence for the alpine case of 19 August 2022
OUTC + 24hours.



PC+1



PC+2

Conclusions and advertisements

- ❑ The modification is available in cycle CY49t1 under namelist option `NVDVAR=5`, thanks to Fabrice Voitus and Karim Yessad.
- ❑ The new formulations may help to further reduce the non-linear residual of the ICI time scheme and to get rid of the noise coming from steep orography, especially in high resolutions.

Tack för din uppmärksamhet!

