Dynamics & Coupling 2005-2006 progress report

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CHMI



Work of: M. Vörös (Hu), R. Brožková (Cz) and F. Váňa (Cz)

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			NH	vertical ve	N	LNH F city [m/	LOV s], N	V STEP	+05	00	
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Fig. 8: NH sl2tl, d_3 + new BBC, diffusion.

BBC for the term $\frac{\partial \tilde{p}}{\partial \pi}$

$$\left(\frac{\partial \tilde{p}}{\partial \pi}\right)_{S} = \frac{\left[-\frac{RT}{p}\nabla p - \nabla \phi + \mathcal{V}\right]_{S} \cdot \nabla \phi_{S} + J_{S} - g\mathcal{W}_{S}}{g^{2} + \left(\nabla \phi_{S}\right)^{2}}$$

$$J_S = \frac{\partial^2 \phi_S}{\partial x^2} u_S^2 + 2 \frac{\partial^2 \phi_S}{\partial x \partial y} u_S v_S + \frac{\partial^2 \phi_S}{\partial y^2} v_S^2$$



BBC for the term $\frac{\partial \tilde{p}}{\partial \pi}$

$$\left(\frac{\partial \tilde{p}}{\partial \pi}\right)_{S} = \frac{\left[-\frac{RT}{p}\nabla p - \nabla \phi + \mathcal{V}\right]_{S} \cdot \nabla \phi_{S} + J_{S} - g\mathcal{W}_{S}}{g^{2} + \left(\nabla \phi_{S}\right)^{2}}$$
$$J_{S} = \frac{\partial^{2}\phi_{S}}{\partial x^{2}}u_{S}^{2} + 2\frac{\partial^{2}\phi_{S}}{\partial x\partial y}u_{S}v_{S} + \frac{\partial^{2}\phi_{S}}{\partial y^{2}}v_{S}^{2}$$

Here
$$\mathcal{V}$$
 and \mathcal{W} are the source terms of momentum containing coriolis, diabatic tendencies and horizontal diffusion.





Possible alternatives:





Possible alternatives:

Switch off HD





Possible alternatives:

- Switch off HD
- Introduce extra spectral computation



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- Use HD computed in GP space



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Semi-Lagrangian Horizontal Diffusion = way to control damping properties of the SL interpolation according to the flow deformation.

 \Rightarrow SLHD solution for ALADIN NH dynamics?



ALADIN 2TL SISL:

$$X_F^+ = \left(1 - \frac{\Delta t}{2}\mathcal{L}\right)^{-1} \left[\underbrace{\left(1 + \frac{\Delta t}{2}\mathcal{L}\right)X_O^- + \Delta t\mathcal{F}_O^- + \frac{\Delta t}{2}\mathcal{N}_O^* + \frac{\Delta t}{2}\mathcal{N}_F^*}_{SLHD}\right]$$

Extra spectral diffusion (= supporting diffusion) is needed for u, v and d (having ϕ_S in their \mathcal{N}).



Reference experiment #1 NH vertical velocity [m/s], NSTEP = +0500

A recreation of a 2D experiment of Jan Masek Nonhydrostatic, nonlinear, Bell shaped mountain Using diffusion - expecting a chimney



Spectral diffusion

∇ 28th EWGLAM + 13th SRNWP Meetings, 9th - 12th of October, Zurich – p.5

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Reference experiment cy29 SLHD #1 NH vertical velocity [m/s], NSTEP = +0500

A recreation of	a 2D experiment o	f Jan Masek						
Nonhydrostatic,	Nonhydrostatic, nonlinear, Bell shaped mountain							
Using SLHD, normal diffusion strength, DiagBBC								
LNHDYN=.T.	LTWOTL=.T.	NSITER=3						
LPC_FULL=.T.	LPC_NESC=.T.	LPC_OLD=.F.						
LADVF=.F.	LGWADV=.F.	LRDBBC=.T.						
RRDXTAU=551.1352	RDAMPDIVS=1.	RDAMPVORS=5.						
SIPR=90000.	SITR=300.	SITRA=50.						
NVDVAR=3	NPDVAR=2	ND4SYS=1						
REPONBT=20000.	REPONTAU=100.	REPONTP=29500.						
NSPONGE=2	LSLHD_W=.T.	LSLHD_SVD=.T.						
SLHDA0=0.25	SLHDB=4.	SLHDD00=6.5E-5	ZSLHDP1=1.7					
ZSLHDP3=0.6	ALPHINT=0.15	GAMMAX0=0.15	SLHDKMAX n/a					
RDAMPVORS=5.	RDAMPDIVS=1.	RDAMPVDS n/a	REXPDHS=6.					
SLEVDHS=1.	SLEVDHS2 n/a	gild point						
			400 400					



GMT 2005 Oct 6 14:30:31 experiment: R016



- Spectral diffusion
- Default SLHD

Reference experiment cy29 SLHD #2 NH vertical velocity [m/s], NSTEP = +0500

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A rec	A recreation of a 2D experiment of Jan Masek							
Nonhy	Nonhydrostatic, nonlinear, Bell shaped mountain							
Using SLHD, no residual diffusion, diagBBC								
LANDYN=, T. LTWOTL=, T. NSITER=3								
LPC_FULL=.T. LPC_NESC=.T. LPC_OLD=.F.								
LADVF=.F. LGWADV=.F. LKDBBC=.T.								
KRDAIAU=U. KDAMPDIVS=1. KDAMPVUKS=5.								
DIPR-	90000. D-2	SIIK-	.300.	STIRA-	-50.			
DEDON	R-3 R-30000		TTATI-100	ND4513				
NCDON	CE-2	J. KEFON	W- T	TGIUD	200- T			
ST.HDA	0=0 25	ST.HDE	w=.1.	ST'HUDU	0=6 5E-5	7.ST.HDD1=	1 7	
ZSLHD	P3=0 6	AL.PHI	NT=0 15	GAMMAX	0=0.35	SLHDKMAX	'1./a	
RDAMP	VORS=5.	RDAME	DTVS=1.	RDAMPV	DS n/a	REXPDHS=	, u	
SLEVD	HS=1.	SLEVE	HS2 n/a	SDRED=	1.			
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GMT 2	006 Jan 314:2	1:34 experiment:	R018				sten:	J.002

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CE



- Default SLHD
- Only GP part of SLHD

Reference experiment cy29 SLHD - Tuned RDAMPVDS NH vertical velocity [m/s], NSTEP = +0500

Nonhydrostatic,	nonlinear, Bell sl	haped mountain				
Using SLHD, normal diffusion strength, DiagBBC						
LNHDYN=.T.	LTWOTL=.T.	NSITER=3				
LPC_FULL=.T.	LPC_NESC=.T.	LPC_OLD=.F.				
LADVF=.F.	LGWADV=.F.	LRDBBC=.T.				
RRDXTAU=551.1352	RDAMPDIVS=1.	RDAMPVORS=5.				
SIPR=90000.	SITR=300.	SITRA=50.				
NVDVAR=3	NPDVAR=2	ND4SYS=1				
REPONBT=20000.	REPONTAU=100.	REPONTP=29500.				
NSPONGE=2	LSLHD_W=.T.	LSLHD_SVD=.T.				
SLHDA0=0.25	SLHDB=4.	SLHDD00=6.5E-5	ZSLHDP1=1.7			
ZSLHDP3=0.6	ALPHINT=0.15	GAMMAX0=0.15	SLHDKMAX n/a			
RDAMPVORS=5.	RDAMPDIVS=1.	RDAMPVDS=15.	REXPDHS=6.			
SLEVDHS=1.	SLEVDHS2 n/a	SDRED=1.				



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CE

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GMT 2006 Apr 3 10:00:04 experiment: R060

- Spectral diffusion
- Default SLHD
- Only GP part of SLHD
- New SLHD

AROME instantaneous rainfalls 19-Oct-2005 00 UTC











AROME MIDPYR LRDBBC+SLHD

Work of: J. Vivoda (Sk)

 $T^*\text{, }T^*_a\Rightarrow T^*(\eta)\text{, }T^*_a(\eta)$





Work of: J. Vivoda (Sk)

 $T^*\text{, }T^*_a\Rightarrow T^*(\eta)\text{, }T^*_a(\eta)$

System becomes more complicated:

- Non trivial setup for $T^*(\eta)$, $T^*_a(\eta)$
- No analyze for optimal $T^*(\eta)$, $T^*_a(\eta)$ setting (at the moment)
- Helmholtz solver becomes the two equations system



2D explicit convection test

Explicit Convection Experiment perturbation of potential temperature [K], NSTEP = +0000



Initial state





 ∇ 28th EWGLAM + 13th SRNWP Meetings, 9th - 12th of October, Zurich – p.8

2D explicit convection test



$\Delta t = 0.5 \text{ s} (\Delta x = \Delta z = 100 \text{ m})$





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2D explicit convection test



$\Delta t = 20$. s ($\Delta x = \Delta z = 100$ m)



FACE

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Work of: J. Vivoda (Sk)

 VFE scheme successfully implemented into the HY model (Untch and Hortal)





- VFE scheme successfully implemented into the HY model (Untch and Hortal)
- Is it extensible to the NH dynamics? (Bénard compatibility and Vivoda - stability)



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- Is it extensible to the NH dynamics? (Bénard compatibility and Vivoda - stability)
- The only non-local operations in the vertical are integrations in HY dynamics (SL version). In NH dynamics also derivatives plays important role (structure equation contains vertical laplacian).



- VFE scheme successfully implemented into the HY model (Untch and Hortal)
- Is it extensible to the NH dynamics? (Bénard compatibility and Vivoda - stability)
- The only non-local operations in the vertical are integrations in HY dynamics (SL version). In NH dynamics also derivatives plays important role (structure equation contains vertical laplacian).
- First version of VFE implemented to the code \rightarrow stable, efficient (2-3 % extra CPU) but (for the moment) noisy.



- VFE scheme successfully implemented into the HY model (Untch and Hortal)
- Is it extensible to the NH dynamics? (Bénard compatibility and Vivoda - stability)
- The only non-local operations in the vertical are integrations in HY dynamics (SL version). In NH dynamics also derivatives plays important role (structure equation contains vertical laplacian).
- First version of VFE implemented to the code \rightarrow stable, efficient (2-3 % extra CPU) but (for the moment) noisy.
- Plan to code a hybrid FE/FD system with interchangeable parts.





NLNH02 test

perturbation of V-wind [m/s], NSTEP = +0500

TSTEP test: 5 2TL ICI NESC scheme NSITER=1

=TRUE

=FALSE

= 3

= 3

LVERTFE

LVFE GW FD

NVSCH

NVDER

LVFE_LAPL_FD =FALSE

LVFE_UVH_FD =FALSE



FD scheme versus full VFE





NLNH02 test

perturbation of V-wind [m/s], NSTEP = +0500

TSTEP test: 5 2TL ICI NESC scheme NSITER=1

=FALSE

=3

=3

LVFE_LAPL_FD =FALSE

LVFE_UVH_FD =FALSE

LVFE GW FD =FALSE

LVERTFE

NVSCH

NVDER

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NLNH02 test

perturbation of V-wind [m/s], NSTEP = +0500

TSTEP test: 5 2TL ICI NESC scheme NSITER=1

=TRUE

= 3

LVERTFE

NVSCH

LVFE_LAPL_FD =TRUE

LVFE_UVH_FD =TRUE

LVFE GW FD =TRUE

LVFE_LAPL_BC_FD =TRUE

NLNH02 test

perturbation of V-wind [m/s], NSTEP = +0500

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TSTEP test: 5 2TL ICI NESC scheme NSITER=1

=FALSE

=3

=3

LVFE_LAPL_FD =FALSE

LVFE_UVH_FD =FALSE

LVFE GW FD =FALSE

LVERTFE

NVSCH

NVDER



FD scheme versus FD with VFE integral operators

NLNH02 test

perturbation of V-wind [m/s], NSTEP = +0500

TSTEP test: 5 2TL ICI NESC scheme NSITER=1

=TRUE

LVFE_LAPL_FD =FALSE

LVFE_UVH_FD =TRUE

LVFE GW FD =TRUE

LVFE_LAPL_BC_FD =TRUE

LVERTFE

NLNH02 test

perturbation of V-wind [m/s], NSTEP = +0500

TSTEP test: 5 2TL ICI NESC scheme NSITER=1

=FALSE

=3

LVERTFE

NVSCH

LVFE_LAPL_FD =FALSE

LVFE_UVH_FD =FALSE

LVFE GW FD =FALSE



FD scheme vs. FD with VFE integ. and laplacian oper.

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Work of: J. Mašek (Sk) and F. Váňa (Cz)

Motivation: SLHD affects conservative properties of the model \Rightarrow need to an improvement of the SL interpolators accuracy.



MSL pressure RMSE and BIAS for 15 days of parallel run





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Work of: J. Mašek (Sk) and F. Váňa (Cz)

Motivation: Performance of the local splines is not superior to the Lagrangian cubic interpolation in SL.



temperature RMSE and BIAS for 15 days of parallel run





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WARM + COLD BUBBLE TEST perturbation of potential temperature [K], NSTEP = +0120

init_102_wcb2_eta, eta-coordinate master al29t2mxl 02 sx6, (A1, A2) = (-1/3, 1/2), .NOT.LOM NH sl2tl, (NPDVAR, NVDVAR) = (2, 3), NSITER = 1, LPC_FULL, LPC_NESC, LGWADV .NOT.LQM[x], .NOT.LQMH[x], LRSPLINE_[x], N[x]LAG = 3 TSTEP = 5.0 s DELY = 10 m DELZ = 10 m P00 = 101325 Pa THETA00 = 300 K SIPR = 90000 Pa SITR = 350 K SITRA = 100 K RRDXTAU = 0





Bubble test, after 10 minutes

• Lagrangian cubic

WARM + COLD BUBBLE TEST perturbation of potential temperature [K], NSTEP = +0120

init_102_wcb2_eta, eta-coordinate master al29t2mxl 02 sx6, (A1, A2) = (-7/15, 4/5), .NOT.LOM NH sl2tl, (NPDVAR, NVDVAR) = (2, 3), NSITER = 1, LPC_FULL, LPC_NESC, LGWADV .NOT.LQM[x], .NOT.LQMH[x], LRSPLINE_[x], N[x]LAG = 3 TSTEP = 5.0 s DELY 10 m = DELZ = 10 m P00 = 101325 Pa THETA00 = 300 K SIPR = 90000 Pa SITR = 350 K SITRA = 100 K RRDXTAU = 0





Bubble test, after 10 minutes

• Lagrangian cubic

Splines

WARM + COLD BUBBLE TEST perturbation of potential temperature [K], NSTEP = +0600

init_102_wcb2_eta, eta-coordinate master al29t2mxl 02 sx6 NH euler, (NPDVAR, NVDVAR) = (2, 3), NSITER = 1, LPC_OLD TSTEP = 1.0 s DELY 10 m DELZ = 10 m = P00 = 101325 Pa THETA00 = 300 K = 90000 Pa SITR = 250 K SITRA = 250 K STPR RRDXTAU = 0





Bubble test, after 10 minutes

- Lagrangian cubic
- Splines
- Eulerian adv.

WARM + COLD BUBBLE TEST perturbation of potential temperature [K], NSTEP = +0120

init_102_wcb2_eta, eta-coordinate master al29t2mxl 02 sx6, (A1, A2) = (0, 0), .NOT.LOM NH sl2tl, (NPDVAR, NVDVAR) = (2, 3), NSITER = 1, LPC_FULL, LPC_NESC, LGWADV .NOT.LQM[x], .NOT.LQMH[x], LRSPLINE_[x], N[x]LAG = 3 TSTEP = 5.0 s DELY = 10 m DELZ = 10 m = 101325 Pa THETA00 = 300 K P00 SIPR = 90000 Pa SITR = 350 K SITRA = 100 K RRDXTAU = 0





Bubble test, after 10 minutes

- Lagrangian cubic
- Splines
- Eulerian adv.
- Linear

Family of two parametric cubic interpolators

$$\begin{split} \mathbf{F}(\mathbf{x},\mathbf{y}) &= \mathbf{w_0}(\mathbf{x})\mathbf{y_0} + \mathbf{w_1}(\mathbf{x})\mathbf{y_1} \\ &+ \mathbf{w_1}(1-\mathbf{x})\mathbf{y_2} + \mathbf{w_0}(1-\mathbf{x})\mathbf{y_3} \end{split}$$

where

$$\begin{array}{lll} w_0(x) &=& a_1x + a_2x^2 - (a_1 + a_2)x^3 \\ w_1(x) &=& 1 + (a_2 - 1)x - (3a_1 + 4a_2)x^2 + 3(a_1 + a_2)x^3 \end{array}$$



Dimensionless damping rate

Damping factor for N = 100, m = 10

Damping factor for N = 100, m = 40









TL/AD of the ALADIN SL

Work of: F. Váňa (Cz)





TL/AD of the ALADIN SL

Work of: F. Váňa (Cz)

Convergence for the TL code:

	$\lim_{\epsilon \to 0} \frac{M(x + \epsilon \delta x) - M}{\mathcal{M}'(\epsilon \delta x)}$	$\frac{\epsilon(x)}{2}, \epsilon = \epsilon_0 10^{\lambda}$
	Eulerian advection Δt =120s	SL advection Δt =450s
$\lambda = 0$	RAT = 0.4922389696335498E+00	RAT =7944390364435365E+01
λ = -1	RAT = 0.9500193013364470E+00	RAT =4770497575992165E+00
λ = -2	RAT = <mark>0.99</mark> 50083001890732E+00	RAT = 0.6874108246125125E+00
λ = -3	RAT = 0.9995037024689268E+00	RAT = 0.9601433242017338E+00
$\lambda = -4$	RAT = <mark>0.9999</mark> 513959612562E+00	RAT = 0.9943026809878674E+00
λ = -5	RAT = <mark>0.1000</mark> 315146923774E+01	RAT = 0.9999531009073782E+00
λ = -6	RAT = 0.1001714189087304E+01	RAT = 0.1001665349367836E+01
λ = -7	RAT = 0.1007310357741422E+01	RAT = 0.1027349076274704E+01
λ = -8	RAT = 0.1119233730823803E+01	RAT = 0.8561242302289194E+00
λ = -9	RAT = 0.5596168654119013E+01	RAT = 0.4280621151144597E+01
λ = -10	RAT = 0.00000000000000E+00	RAT = 0.00000000000000E+00
_		





TL/AD of the ALADIN SL

Work of: F. Váňa (Cz)

Convergence for the TL code:

	$\lim_{\epsilon \to 0} \frac{M(x + \epsilon \delta x) - M}{\mathcal{M}'(\epsilon \delta x)}$	$\frac{(x)}{2}, \epsilon = \epsilon_0 10^{\lambda}$
	Eulerian advection Δt =120s	SL advection Δt =450s
$\lambda = 0$	RAT = 0.9685219082957116E+00	RAT = 0.1094034387101322E+01
λ = -1	RAT = 0.9970618603595810E+00	RAT = 0.1008012195504008E+01
λ = -2	RAT = 0.9997073040468342E+00	RAT = 0.1002141025110223E+01
λ = -3	RAT = 0.9999707398884352E+00	RAT = 0.1000160788422592E+01
$\lambda = -4$	RAT = <mark>0.99999</mark> 70679271253E+00	RAT = 0.1000099605664519E+01
λ = -5	RAT = 0.9999995490240665E+00	RAT = 0.1000001139215519E+01
λ = -6	RAT = <mark>0.99999</mark> 87045356886E+00	RAT = 0.1000001847670018E+01
λ = -7	RAT = 0.9999936488857756E+00	RAT = 0.1000041939684409E+01
$\lambda = -8$	RAT = 0.9999533728917936E+00	RAT = 0.1000246087384355E+01
$\lambda = -9$	RAT = 0.9991377690586460E+00	RAT = 0.9994838411148169E+00
λ = -10	RAT = 0.9970808134568164E+00	RAT = 0.1032182685987080E+01





Work of: R. Hamdi (Be) and P. Termonia (Be)

 The way in which the physics is coupled to the dynamics has an influence on the stability and the accuracy



Work of: R. Hamdi (Be) and P. Termonia (Be)

- The way in which the physics is coupled to the dynamics has an influence on the stability and the accuracy
- Simple 1d model simulations using the framework proposed by Staniforth, Wood, Côté (2002) extended in a way to take into account the spectral nature of the models and the difference between the real atmosphere and the background of the linearisation.



Work of: R. Hamdi (Be) and P. Termonia (Be)

- The way in which the physics is coupled to the dynamics has an influence on the stability and the accuracy
- Simple 1d model simulations using the framework proposed by Staniforth, Wood, Côté (2002) extended in a way to take into account the spectral nature of the models and the difference between the real atmosphere and the background of the linearisation.
- This study was restricted to explicit, semi implicit and implicit physics parameterizations (over-implicitness not treated).



Possibilities to organize a time step

coupling of the physics parameterization <u>before</u> or <u>after</u> the explicit part of the dynamics



Possibilities to organize a time step

- coupling of the physics parameterization <u>before</u> or <u>after</u> the explicit part of the dynamics
- coupling of the physics to the dynamics at different positions (in space and time with respect to dyn.) on the SL trajectory



Possibilities to organize a time step

- coupling of the physics parameterization <u>before</u> or <u>after</u> the explicit part of the dynamics
- coupling of the physics to the dynamics at different positions (in space and time with respect to dyn.) on the SL trajectory
- computing the physics parameterization in a parallel or a fractional manner



Possibilities to organize a time step

- coupling of the physics parameterization <u>before</u> or <u>after</u> the explicit part of the dynamics
- coupling of the physics to the dynamics at different positions (in space and time with respect to dyn.) on the SL trajectory
- computing the physics parameterization in a parallel or a <u>fractional</u> manner
- coupling the physics to the dynamics by updating the model state and using this for the dynamics, or computing the physics tendency and the dynamics tendencies separately and adding them to get the update, in other words to treat the physics/dynamics in a fractional or a sequential manner



A/A/A vs. SLAVEPP

	A/A/A	SLAVEPP
phys. before/after dyn.	before	computed after and averaged
on SL traj.	at t	at $t + \Delta t$
parallel / sequential physics calls	parallel	sequential
parallel /sequential phys dyn. coupling	sequential	parallel



Results

Always couple the physics to the air parcel along the SL trajectory. Otherwise the properties (stability and accuracy) depend on the advection.



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- The structure of A/A/A is more stable but less accurate in case of a single diffuse process.



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- The structure of A/A/A is more stable but less accurate in case of a single diffuse process.
- In parallel physics coupling one should couple the diffusive processes last. (This confirms results in Dubal et al. (2004)).



- Always couple the physics to the air parcel along the SL trajectory. Otherwise the properties (stability and accuracy) depend on the advection.
- The structure of A/A/A is more stable but less accurate in case of a single diffuse process.
- In parallel physics coupling one should couple the diffusive processes last. (This confirms results in Dubal et al. (2004)).
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- In the A/A/A framework, coupling the physics after the dynamics gives a more accurate treatment of the steady state. This might be beneficial for climate simulations.
- If the physics were treated in a semi-implicit way (in the A/A/A context) we would have the same stability and also second-order accuracy as in the SLAVEPP approach. This is maybe not practical but nevertheless a nice surprise because it means that (in the 1d model) one could get the same benefits of the time step reorganization, by an internal reorganization of the physics.

Outcome for A/A/A

If both forcing and diffusive processes are present, a SLAVEPP kind of time step becomes superior to the A/A/A one *IF* the diffusive processes are coupled *LAST*. This will also lead to a more accurate steady state and less climate drift.



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- There seems to be no better option with respect of phys-dyn coupling to increase the existing stability of the physics in the A/A/A framework.
- Publication with the detailed guidelines is in preparation (manuscript can be obtained from *piet.termonia@oma.be*)



Work of: F. Voitus (Fr) and P. Termonia (Be)

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- This is related to a deeper question: "Can one impose LBCs independently from the details of the dynamics integration scheme?"
- In a spectral model we are forced to address this question, but in a gridpoint model maintained in a huge collaboration with different kinds of researchers working together, this will have to be addressed too.

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- Proposal: rely on sub-stepping



Is it feasible?

Tests with a leapfrog scheme with following space-time structure





First test in Shallow water model

Inside the domain (crosses): SISL 2TL

Near the boundary to compute 3 points at $t + \Delta t$ (solid dots): leapfrog Asselin or leapfrog trapezoidal by sub-stepping with interval τ



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- The problem is being replaced to data flow...: We need an extra large stencil near the boundary (solid dots at t).

